

Yang Z, Rodríguez CE. 2013. Searching for efficient Markov chain Monte Carlo proposal kernels. *Proc Natl Acad Sci USA* 110:19307–19312

SI Text S3 (extended version).  $P_{\text{jump}}$  for uniform and Bactrian kernels on  $N(0, 1)$  target

Ziheng Yang, 12 November 2013

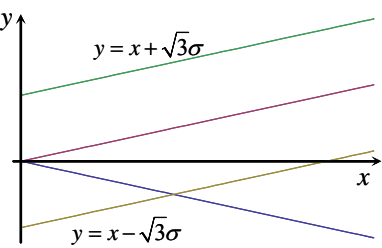
(The equation numbers here differ from those in the paper)

Note that both uniform and Bactrian kernels are symmetrical with  $q(y|x) = q(x|y) = q(|y - x|)$ . Thus

$$\begin{aligned} P_{\text{jump}} &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \phi(x)q(y|x)I_{\phi(y) > \phi(x)} dx dy + \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \phi(x)q(y|x)I_{\phi(y) \leq \phi(x)} \frac{\phi(y)}{\phi(x)} dx dy \\ &= 2 \iint_{\phi(y) > \phi(x)} \phi(x)q(y|x) dy dx && \leftarrow \text{detailed balance} \\ &= 4 \int_0^{\infty} \int_{-x}^x \phi(x)q(y|x) dy dx. && \leftarrow \text{symmetry} \end{aligned} \quad (1)$$

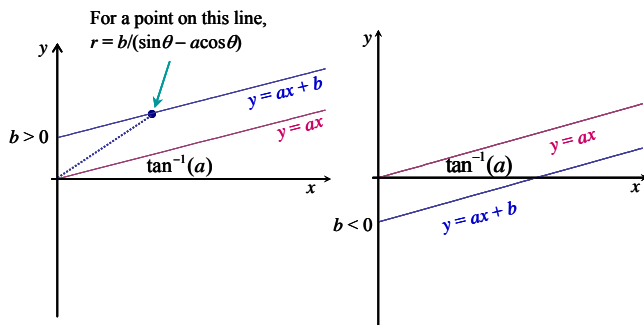
The mass of  $\phi(x)q(y|x)$  left of the  $y$  axis is the same as right of the  $y$  axis. Also right of the  $y$  axis,  $\phi(y) > \phi(x)$  iff  $-x < y < x$ .

For the uniform kernel, we have



$$\begin{aligned} P_{\text{jump}} &= \frac{2}{\sqrt{3}\sigma} \left[ \int_0^{\frac{\sqrt{3}\sigma}{2}} \int_{-x}^x \phi(x) dy dx + \int_{\frac{\sqrt{3}\sigma}{2}}^{\infty} \int_{x-\sqrt{3}\sigma}^x \phi(x) dy dx \right] \\ &= \frac{2}{\sqrt{3}\sigma} \left[ 2 \int_0^{\frac{\sqrt{3}\sigma}{2}} x\phi(x) dx + \sqrt{3}\sigma \int_{\frac{\sqrt{3}\sigma}{2}}^{\infty} \phi(x) dx \right] \\ &= \frac{2}{\sqrt{3}\sigma} \left[ \sqrt{\frac{2}{\pi}} \left( 1 - e^{-\frac{3\sigma^2}{8}} \right) + \sqrt{3}\sigma \left( 1 - \Phi\left(\frac{\sqrt{3}\sigma}{2}\right) \right) \right] \\ &= \sqrt{\frac{8}{3\pi\sigma^2}} \left( 1 - e^{-\frac{3\sigma^2}{8}} \right) + 2 \left( 1 - \Phi\left(\frac{\sqrt{3}\sigma}{2}\right) \right). \end{aligned}$$

Before we consider the Bactrian kernel, we give the following integral, which we need later.



$$\begin{aligned} I(a, b) &= \int_0^{\infty} \phi(x)\Phi(ax+b) dx = \int_0^{\infty} \int_{-\infty}^{ax+b} \frac{1}{2\pi} \exp\left\{-\frac{x^2+y^2}{2}\right\} dy dx, \\ &-\infty < a, b < \infty. \text{ As } I(-a, b) = \frac{1}{2} - I(a, -b), \text{ we} \\ &\text{consider the case of } a > 0 \text{ only. We change to the} \\ &\text{polar system. If } b \geq 0, \end{aligned}$$

$$\begin{aligned} I_+(a, b) &= \frac{1}{2\pi} \int_{-\frac{\pi}{2}}^{\tan^{-1}(a)} \int_0^{\infty} e^{-\frac{r^2}{2}} r dr d\theta + \frac{1}{2\pi} \int_{\tan^{-1}(a)}^{\frac{\pi}{2}} \int_0^{\frac{b}{\sin\theta - a\cos\theta}} e^{-\frac{r^2}{2}} r dr d\theta \\ &= \frac{1}{2\pi} \int_{-\frac{\pi}{2}}^{\tan^{-1}(a)} 1 d\theta + \frac{1}{2\pi} \int_{\tan^{-1}(a)}^{\frac{\pi}{2}} \left[ 1 - \exp\left\{-\frac{b^2}{2(\sin\theta - a\cos\theta)^2}\right\} \right] d\theta \\ &= \frac{1}{2\pi} \left[ \tan^{-1}(a) + \frac{\pi}{2} \right] + \frac{1}{2\pi} \left[ \frac{\pi}{2} - \tan^{-1}(a) - \int_{\tan^{-1}(a)}^{\frac{\pi}{2}} \exp\left\{-\frac{b^2}{2(\sin\theta - a\cos\theta)^2}\right\} d\theta \right] \\ &= \frac{1}{2} - \frac{1}{2\pi} \int_{\tan^{-1}(a)}^{\frac{\pi}{2}} \exp\left\{-\frac{b^2}{2(\sin\theta - a\cos\theta)^2}\right\} d\theta. \end{aligned} \quad (4)$$

Similarly if  $b \leq 0$ ,

$$I_-(a, b) = \frac{1}{2\pi} \int_{-\frac{\pi}{2}}^{\tan^{-1}(a)} \int_{\frac{b}{\sin\theta - a \cos\theta}}^{\infty} e^{-\frac{r^2}{2}} r dr d\theta = \frac{1}{2\pi} \int_{-\frac{\pi}{2}}^{\tan^{-1}(a)} \exp\left\{-\frac{b^2}{2(\sin\theta - a \cos\theta)^2}\right\} d\theta. \quad (5)$$

If  $b = 0$ , both  $I_+(a, b)$  and  $I_-(a, b)$  reduce to

$$\int_0^{\infty} \phi(x) \Phi(ax) dx = \frac{1}{4} + \frac{1}{2\pi} \tan^{-1}(a), \text{ with } a > 0. \quad (6)$$

Note that the proposal density for the Bactrian kernel (eq. 7 in the paper) can be written as

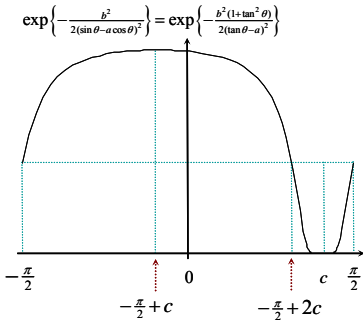
$$q(y|x) = \frac{1}{2} \phi(y; x + m\sigma, (1 - m^2)\sigma^2) + \frac{1}{2} \phi(y; x - m\sigma, (1 - m^2)\sigma^2). \quad (7)$$

Thus from eq. (1),

$$\begin{aligned} P_{\text{jump}} &= 2 \int_0^{\infty} \int_{-x}^x \phi(x) \times \left[ \phi(y; x + m\sigma, (1 - m^2)\sigma^2) + \phi(y; x - m\sigma, (1 - m^2)\sigma^2) \right] dy dx \\ &= 2 \int_0^{\infty} \phi(x) \times \left[ \Phi\left(\frac{-m}{\sqrt{1-m^2}}\right) - \Phi\left(\frac{-2x-m\sigma}{\sqrt{1-m^2}\sigma}\right) + \Phi\left(\frac{m}{\sqrt{1-m^2}}\right) - \Phi\left(\frac{-2x+m\sigma}{\sqrt{1-m^2}\sigma}\right) \right] dx \\ &= 1 - 2 \int_0^{\infty} \phi(x) \times \left[ \Phi\left(\frac{-2x-m\sigma}{\sqrt{1-m^2}\sigma}\right) + \Phi\left(\frac{-2x+m\sigma}{\sqrt{1-m^2}\sigma}\right) \right] dx \\ &= -1 + 2 \int_0^{\infty} \phi(x) \left[ \Phi\left(\frac{2x+m\sigma}{\sqrt{1-m^2}\sigma}\right) + \Phi\left(\frac{2x-m\sigma}{\sqrt{1-m^2}\sigma}\right) \right] dx. \end{aligned} \quad (8)$$

Now define  $a = \frac{2}{\sigma\sqrt{1-m^2}} > 0$ ,  $b = \frac{m}{\sqrt{1-m^2}} > 0$ . Eq. (8) then becomes

$$\begin{aligned} P_{\text{jump}} &= -1 + 2[I_+(a, b) + I_-(a, -b)] \\ &= \frac{1}{\pi} \left[ \int_{-\frac{\pi}{2}}^{\tan^{-1}(a)} \exp\left\{-\frac{b^2}{2(\sin\theta - a \cos\theta)^2}\right\} d\theta - \int_{\tan^{-1}(a)}^{\frac{\pi}{2}} \exp\left\{-\frac{b^2}{2(\sin\theta - a \cos\theta)^2}\right\} d\theta \right]. \end{aligned} \quad (9)$$



The integrand is the same in the two integrals and is 0 at  $\theta = c = \tan^{-1}(a)$ , and symmetrical around  $c$ ; in other words it is the same at  $c - d$  and at  $c + d$ , for  $0 < d < \pi/2 - c$ . Thus

$$P_{\text{jump}} = \frac{1}{\pi} \int_{-\frac{\pi}{2}}^{-\frac{\pi}{2} + 2c} \exp\left\{-\frac{b^2}{2(\sin\theta - a \cos\theta)^2}\right\} d\theta. \quad (10)$$

Now the integral is over an interval of length  $2c$ , and the integrand is symmetrical around the center: it is the same at  $(-\pi/2 + c) - d'$  and at  $(-\pi/2 + c) + d'$  for  $0 < d' < c$ . Thus

$$P_{\text{jump}} = \frac{2}{\pi} \int_{-\frac{\pi}{2}}^{-\frac{\pi}{2} + \tan^{-1}(a)} \exp\left\{-\frac{b^2}{2(\sin\theta - a \cos\theta)^2}\right\} d\theta = \frac{2}{\pi} \int_0^{\tan^{-1}(a)} \exp\left\{-\frac{b^2}{2(\cos\theta + a \sin\theta)^2}\right\} d\theta \quad (11)$$

Let  $t = \tan\theta$ , with  $d\theta = 1/(1 + t^2) dt$ . Then

$$P_{\text{jump}} = \frac{2}{\pi} \int_0^a \frac{1}{1 + t^2} \exp\left\{-\frac{b^2(1 + t^2)}{2(1 + at)^2}\right\} dt. \quad (12)$$

This reduces to eq. (11) in the paper if  $b = 0$  (if  $m = 0$ ).